

Completeness results for the ZX-calculus

Miriam Backens

Department of Computer Science, University of Oxford

Quantum Physics and Logic, Kyoto, 2014

Outline

Motivation

Approximate completeness for single qubits

Completeness for stabilizer diagrams

Conclusions & Outlook

Motivation

- ▶ ZX-calculus is incomplete, even for single qubits; not obvious how to complete it
- ▶ instead: look for fragments of the general calculus which are complete
- ▶ approximate universality: small sets of operators suffice to approximate arbitrary unitaries to any accuracy
- ▶ *approximate completeness*: completeness for such a set

Outline

Motivation

Approximate completeness for single qubits

Completeness for stabilizer diagrams

Conclusions & Outlook

The single-qubit Clifford+T group

An approximately universal group, generated by:

- ▶ single-qubit Clifford group $\mathcal{C}_1 = \langle S, H \rangle$, where

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix} \quad \text{⊗ } \pi/2$$

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \quad \text{⊗ } H$$

- ▶ T gate

$$T = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\pi/4} \end{pmatrix} \quad \text{⊗ } \pi/4$$

The ZX-calculus for the single-qubit Clifford+T group

Generated by $\bullet \pi/2$, \square , and $\bullet \pi/4$ —or $\bullet \pi/4$ and $\bullet \pi/2$

- ▶ *single-qubit*: restrict diagrams to line graphs
- ▶ *Clifford+T group*: restrict phases to multiples of $\pi/4$

(Ignore global phases.)

The ZX-calculus for the single-qubit Clifford+T group

Generated by $\bullet \pi/2$, $\square H$, and $\bullet \pi/4$ —or $\bullet \pi/4$ and $\bullet \pi/2$

- ▶ *single-qubit*: restrict diagrams to line graphs
- ▶ *Clifford+T group*: restrict phases to multiples of $\pi/4$

(Ignore global phases.)

Rules:

$$\begin{array}{c} \bullet \alpha \\ | \\ \bullet \beta \end{array} = \bullet \alpha + \beta$$

$$\begin{array}{c} \bullet \pi \\ | \\ \bullet \alpha \end{array} = \begin{array}{c} \bullet -\alpha \\ | \\ \bullet \pi \end{array}$$

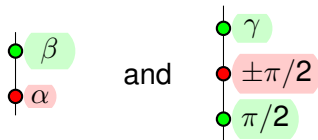
$$\begin{array}{c} \square H \\ | \\ \bullet \alpha \\ | \\ \square H \end{array} = \bullet \alpha$$

$$\square H = \begin{array}{c} \bullet \pi/2 \\ | \\ \bullet \pi/2 \\ | \\ \bullet \pi/2 \end{array}$$

and $\bullet 2n\pi = |$ for integer n

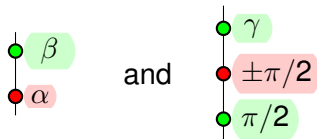
Single-qubit Clifford diagrams

Any single-qubit Clifford diagram can be written uniquely as one of the following, with $\alpha, \beta, \gamma \in \{0, \pi/2, \pi, -\pi/2\}$:

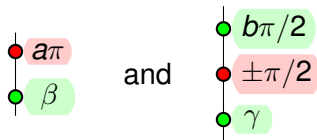


Single-qubit Clifford diagrams

Any single-qubit Clifford diagram can be written uniquely as one of the following, with $\alpha, \beta, \gamma \in \{0, \pi/2, \pi, -\pi/2\}$:



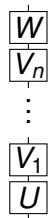
or, equivalently, one of



where $a, b \in \{0, 1\}$ and β, γ as above.

The normal form for single-qubit Clifford+T diagrams

Following [Matsumoto & Amano 2008], any single-qubit Clifford+T diagram is either pure Clifford or it can be written as



where

$$W \in \left\{ \begin{array}{c} | \\ \bullet \pi/2 \\ | \\ \bullet \pi/2 \\ | \end{array} \right\}$$

$$V_k \in \left\{ \begin{array}{cc} \bullet \pi/4 & \bullet 3\pi/4 \\ \bullet \pi/2 & \bullet \pi/2 \end{array} \right\} \text{ for } 1 \leq k \leq n$$

$$U \in \left\{ \begin{array}{cc} \bullet \pi/4 + \alpha & \bullet \pi/4 + \gamma \\ \bullet \beta & \bullet \pm\pi/2 \\ & \bullet \pi/2 \end{array} \right\}$$

with n a non-negative integer and $\alpha, \beta, \gamma \in \{0, \pi/2, \pi, -\pi/2\}$.

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

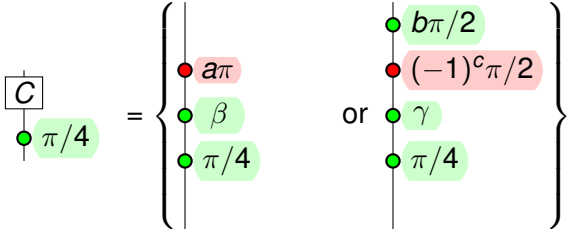
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{l} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $V \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{l} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \quad \begin{array}{l} \bullet b\pi/2 \\ \bullet \pm\pi/2 \\ \bullet \gamma \end{array} \right\}$, then



Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \right\}$, then

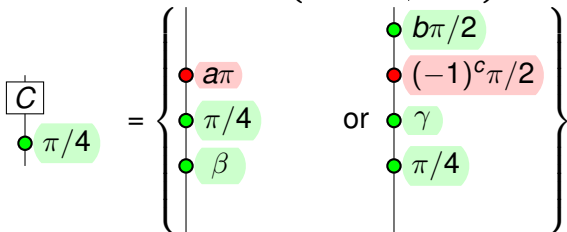
$$\boxed{C} \begin{array}{c} \bullet \pi/4 \end{array} = \left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta + \pi/4 \end{array} \text{ or } \begin{array}{l} \bullet b\pi/2 \\ \bullet (-1)^c \pi/2 \\ \bullet \gamma \\ \bullet \pi/4 \end{array} \right\}$$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{l} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \quad \begin{array}{l} \bullet b\pi/2 \\ \bullet \pm\pi/2 \\ \bullet \gamma \end{array} \right\}$, then



Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{l} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \quad \begin{array}{l} \bullet b\pi/2 \\ \bullet \pm\pi/2 \\ \bullet \gamma \end{array} \right\}$, then

$$\boxed{C} \begin{array}{c} \bullet \pi/4 \end{array} = \left\{ \begin{array}{l} \begin{array}{c} \bullet (-1)^a \pi/4 \\ \bullet a\pi \\ \bullet \beta \end{array} \quad \text{or} \quad \begin{array}{c} \bullet b\pi/2 \\ \bullet (-1)^c \pi/2 \\ \bullet \gamma \\ \bullet \pi/4 \end{array} \end{array} \right\}$$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $V \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \right\}$, then

$$\begin{array}{c} \boxed{C} \\ \bullet \pi/4 \end{array} = \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet -a\pi/2 \\ \bullet a\pi \\ \bullet \beta \end{array} \right\} \text{ or } \left\{ \begin{array}{l} \bullet b\pi/2 \\ \bullet (-1)^c \pi/2 \\ \bullet \gamma \\ \bullet \pi/4 \end{array} \right\}$$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \right\}$, then

$$\boxed{C} \begin{array}{c} \bullet \pi/4 \end{array} = \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet a\pi \\ \bullet \beta + a\pi/2 \end{array} \right\} \text{ or } \left\{ \begin{array}{l} \bullet b\pi/2 \\ \bullet (-1)^c \pi/2 \\ \bullet \gamma \\ \bullet \pi/4 \end{array} \right\}$$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \right\}$, then

$$\boxed{C} \begin{array}{c} \bullet \pi/4 \end{array} = \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet a\pi \\ \bullet \beta + a\pi/2 \end{array} \right\} \text{ or } \left\{ \begin{array}{l} \bullet b\pi/2 \\ \bullet (-1)^c \pi/2 \\ \bullet \pi/4 \\ \bullet \gamma \end{array} \right\}$$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{l} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \quad \begin{array}{l} \bullet b\pi/2 \\ \bullet \pm\pi/2 \\ \bullet \gamma \end{array} \right\}$, then

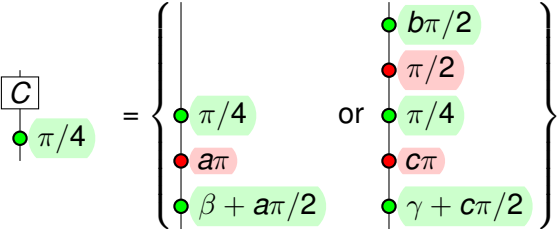
$$\boxed{C} \begin{array}{c} \bullet \pi/4 \end{array} = \left\{ \begin{array}{l} \begin{array}{c} \bullet \pi/4 \\ \bullet a\pi \\ \bullet \beta + a\pi/2 \end{array} \quad \text{or} \quad \begin{array}{c} \bullet b\pi/2 \\ \bullet \pi/2 \\ \bullet c\pi \\ \bullet \pi/4 \\ \bullet \gamma \end{array} \end{array} \right.$$

Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{l} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{l} \bullet a\pi \\ \bullet \beta \end{array} \quad \begin{array}{l} \bullet b\pi/2 \\ \bullet \pm\pi/2 \\ \bullet \gamma \end{array} \right\}$, then

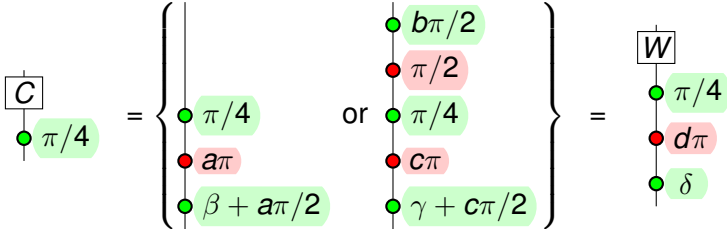


Pushing phases down, part 1

Diagrams are rewritten into normal form by pushing phases towards the bottom of the diagram as much as possible.

- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/4$
- ▶ can push $\bullet \alpha$ and $\bullet \pi$ past $\bullet \pi/2$
- ▶ can push $\bullet \pi$ and $\bullet \pi$ past $\boxed{V} \in \left\{ \begin{array}{c} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \quad \begin{array}{c} \bullet 3\pi/4 \\ \bullet \pi/2 \end{array} \right\}$

Write Clifford operators as $\left\{ \begin{array}{c} \bullet a\pi \\ \bullet \beta \end{array} \quad \begin{array}{c} \bullet b\pi/2 \\ \bullet \pm\pi/2 \\ \bullet \gamma \end{array} \right\}$, then

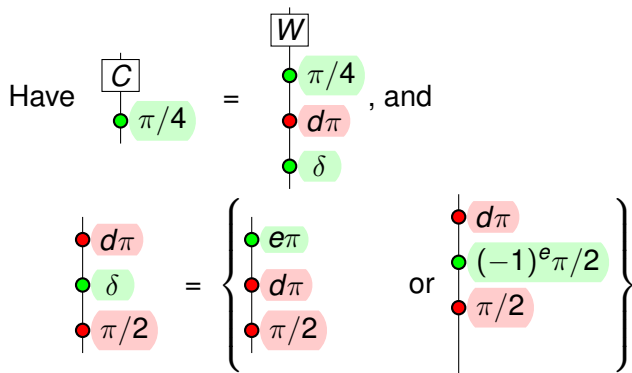


Pushing phases down, part 2

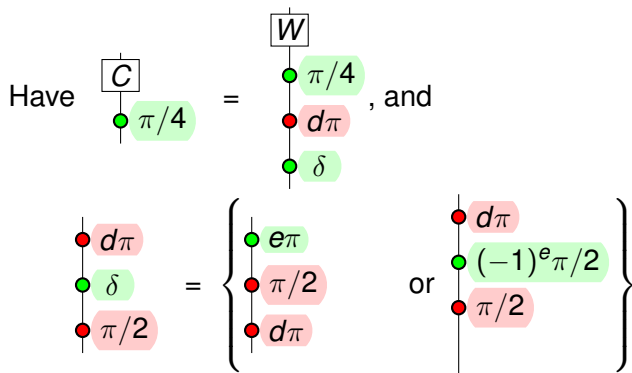
Have \boxed{C} $\begin{array}{c} \bullet \\ \pi/4 \end{array}$ = $\begin{array}{c} \boxed{W} \\ \bullet \\ \pi/4 \\ \bullet \\ d\pi \\ \bullet \\ \delta \end{array}$, and

$$\begin{array}{c} \bullet \\ d\pi \\ \bullet \\ \delta \\ \bullet \\ \pi/2 \end{array} = \left\{ \begin{array}{c} \bullet \\ d\pi \\ \bullet \\ e\pi \\ \bullet \\ \pi/2 \end{array} \right. \text{ or } \left. \begin{array}{c} \bullet \\ d\pi \\ \bullet \\ (-1)^e \pi/2 \\ \bullet \\ \pi/2 \end{array} \right\}$$

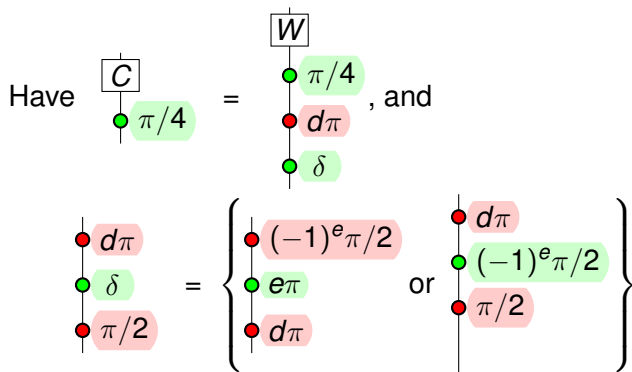
Pushing phases down, part 2



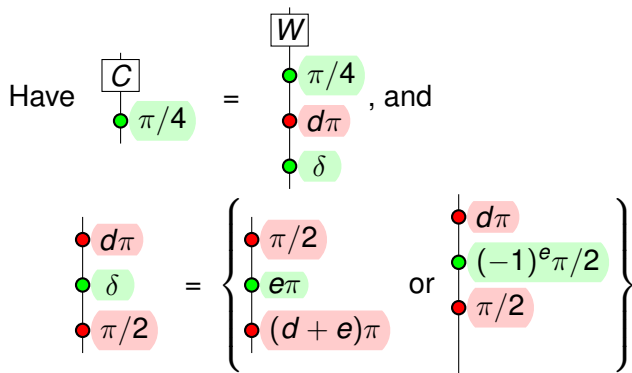
Pushing phases down, part 2



Pushing phases down, part 2



Pushing phases down, part 2



Pushing phases down, part 2

Have \boxed{C} $\begin{array}{c} \bullet \\ \pi/4 \end{array}$ = $\begin{array}{c} \boxed{W} \\ \bullet \\ \pi/4 \\ \bullet \\ d\pi \\ \bullet \\ \delta \end{array}$, and

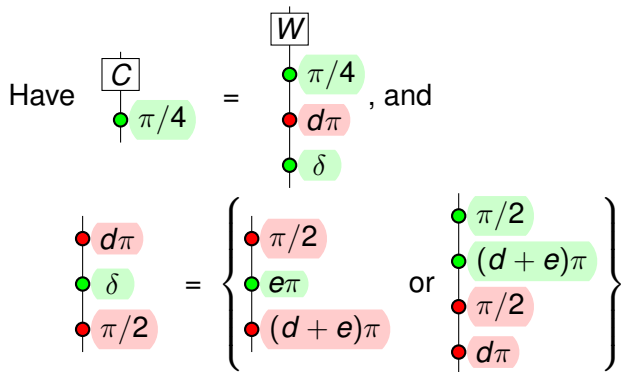
$$\begin{array}{c} \bullet \\ d\pi \\ \bullet \\ \delta \\ \bullet \\ \pi/2 \end{array} = \left\{ \begin{array}{c} \bullet \\ \pi/2 \\ \bullet \\ e\pi \\ \bullet \\ (d+e)\pi \end{array} \right. \text{ or } \left. \begin{array}{c} \bullet \\ (-1)^{d+e}\pi/2 \\ \bullet \\ d\pi \\ \bullet \\ \pi/2 \end{array} \right\}$$

Pushing phases down, part 2


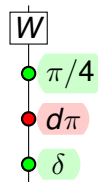
Have \boxed{C} $\begin{array}{c} \bullet \\ \pi/4 \end{array}$ = $\begin{array}{c} \boxed{W} \\ \bullet \\ \pi/4 \\ \bullet \\ d\pi \\ \bullet \\ \delta \end{array}$, and

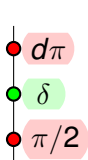
$$\begin{array}{c} \bullet \\ d\pi \\ \bullet \\ \delta \\ \bullet \\ \pi/2 \end{array} = \left\{ \begin{array}{c} \bullet \\ \pi/2 \\ \bullet \\ e\pi \\ \bullet \\ (d+e)\pi \end{array} \right. \text{ or } \left. \begin{array}{c} \bullet \\ (-1)^{d+e}\pi/2 \\ \bullet \\ \pi/2 \\ \bullet \\ d\pi \end{array} \right\}$$

Pushing phases down, part 2

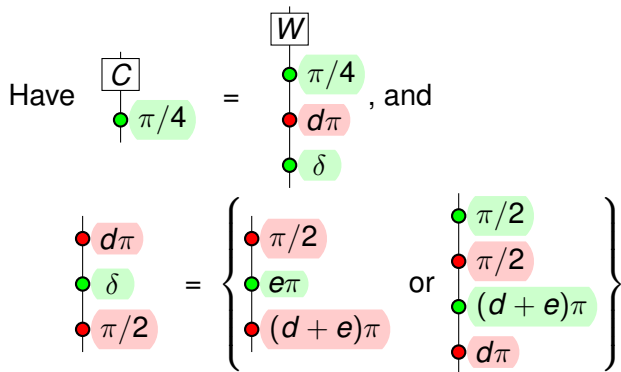


Pushing phases down, part 2

Have \boxed{C}  = \boxed{W} , and

 = $\left\{ \begin{array}{l} \text{red } \pi/2 \\ \text{green } e\pi \\ \text{red } (d+e)\pi \end{array} \right.$ or $\left. \begin{array}{l} \text{green } \pi/2 \\ \text{red } (-1)^{d+e}\pi/2 \\ \text{green } (d+e)\pi \\ \text{red } e\pi \end{array} \right\}$

Pushing phases down, part 2



Pushing phases down, part 2

Have \boxed{C} $\begin{array}{c} \bullet \\ \pi/4 \end{array}$ = $\begin{array}{c} \boxed{W} \\ \bullet \\ \pi/4 \\ \bullet \\ d\pi \\ \bullet \\ \delta \end{array}$, and

$$\begin{array}{c} \bullet \\ d\pi \\ \bullet \\ \delta \\ \bullet \\ \pi/2 \end{array} = \left\{ \begin{array}{c} \bullet \\ \pi/2 \\ \bullet \\ e\pi \\ \bullet \\ (d+e)\pi \end{array} \right. \text{ or } \left. \begin{array}{c} \bullet \\ \pi/2 \\ \bullet \\ \pi/2 \\ \bullet \\ (d+e)\pi \\ \bullet \\ d\pi \end{array} \right\} = \begin{array}{c} \bullet \\ f\pi/2 \\ \bullet \\ \pi/2 \\ \bullet \\ g\pi \\ \bullet \\ h\pi \end{array},$$

Pushing phases down, part 2

Have \boxed{C} $\begin{array}{c} \bullet \\ \pi/4 \end{array}$ = $\begin{array}{c} \boxed{W} \\ \bullet \\ \pi/4 \\ \bullet \\ d\pi \\ \bullet \\ \delta \end{array}$, and

$$\begin{array}{c} \bullet \\ d\pi \\ \bullet \\ \delta \\ \bullet \\ \pi/2 \end{array} = \left\{ \begin{array}{c} \bullet \\ \pi/2 \\ \bullet \\ e\pi \\ \bullet \\ (d+e)\pi \end{array} \text{ or } \begin{array}{c} \bullet \\ \pi/2 \\ \bullet \\ (d+e)\pi \\ \bullet \\ d\pi \end{array} \right\} = \begin{array}{c} \bullet \\ f\pi/2 \\ \bullet \\ \pi/2 \\ \bullet \\ g\pi \\ \bullet \\ h\pi \end{array},$$

therefore

$$\begin{array}{c} \boxed{C} \\ \boxed{V} \end{array} = \begin{array}{c} \boxed{W} \\ \boxed{V'} \\ \bullet \\ d\pi \\ \bullet \\ e\pi \end{array}, \quad \begin{array}{c} \boxed{C} \\ \boxed{U} \end{array} = \begin{array}{c} \boxed{W} \\ \boxed{U'} \end{array}, \quad \text{and} \quad \begin{array}{c} \bullet \\ a\pi \\ \bullet \\ b\pi \\ \boxed{V} \end{array} = \begin{array}{c} \boxed{V'} \\ \bullet \\ e\pi \\ \bullet \\ f\pi \end{array}.$$

Rewriting diagrams into normal form, starting with a Clifford diagram

The single-qubit Clifford+T group is generated by $\bullet \pi/4$ and $\bullet \pi/2$ so it suffices to check what happens when a pure Clifford operator or a normal form diagram is composed with one of these.

Now for any Clifford unitary \boxed{C} ,

$$\boxed{C} \bullet \pi/2 \text{ is pure Clifford, and } \boxed{C} \bullet \pi/4 = \boxed{U} \in \left\{ \begin{array}{l} \bullet \pi/4 + \alpha \\ \bullet \pi/2 \\ \bullet \pi/4 + \gamma \\ \bullet \pm \pi/2 \\ \bullet 3\pi/2 \end{array} \right\}.$$

Rewriting diagrams into normal form, starting with a normal form diagram

For any $\boxed{W} \in \left\{ \begin{array}{l} \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right\}$ and $\boxed{V_n} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \right\}$,

$$\begin{array}{c} \bullet \pi/2 \\ \boxed{W} \\ \boxed{V_n} \end{array} = \boxed{C} = \begin{array}{c} \boxed{W'} \\ \boxed{V'_n} \\ \bullet a\pi \\ \bullet b\pi \end{array}$$

and

$$\begin{array}{c} \bullet \pi/4 \\ \boxed{W} \\ \boxed{V_n} \end{array} = \left\{ \boxed{C'} \text{ or } \begin{array}{c} \boxed{V_{n+1}} \\ \boxed{V_n} \end{array} \right\}$$

Rewriting diagrams into normal form, starting with a normal form diagram

For any $\boxed{W} \in \left\{ \begin{array}{l} \bullet \pi/2 \\ \bullet \pi/2 \end{array} \right\}$ and $\boxed{V_n} \in \left\{ \begin{array}{l} \bullet \pi/4 \\ \bullet \pi/2 \end{array} \right\}$,

$$\begin{array}{c} \bullet \pi/2 \\ \boxed{W} \\ \boxed{V_n} \end{array} = \boxed{C} = \begin{array}{c} \boxed{W'} \\ \boxed{V'_n} \\ \bullet a\pi \\ \bullet b\pi \end{array}$$

and

$$\begin{array}{c} \bullet \pi/4 \\ \boxed{W} \\ \boxed{V_n} \end{array} = \left\{ \boxed{C'} \text{ or } \begin{array}{c} \boxed{V_{n+1}} \\ \boxed{V_n} \end{array} \right\}$$

Case $n = 0$ is similar.

No normal form diagram represents the identity operator

Can write any single-qubit density operator as $xX + yY + zZ$, where $x, y, z \in \mathbb{R}$ and X, Y, Z are the Pauli matrices.

Clifford unitaries act on the vectors (x, y, z) by permuting the elements and adding minus signs; T sends

$$(x, y, z) \mapsto \frac{1}{\sqrt{2}} (x - y, x + y, z\sqrt{2}).$$

No normal form diagram represents the identity operator

Can write any single-qubit density operator as $xX + yY + zZ$, where $x, y, z \in \mathbb{R}$ and X, Y, Z are the Pauli matrices.

Clifford unitaries act on the vectors (x, y, z) by permuting the elements and adding minus signs; T sends

$$(x, y, z) \mapsto \frac{1}{\sqrt{2}} (x - y, x + y, z\sqrt{2}).$$

States resulting from application of a normal form operator to $|0\rangle$ will have vectors of the form

$$\frac{1}{\sqrt{2^m}} (x_1 + x_2\sqrt{2}, y_1 + y_2\sqrt{2}, z_1 + z_2\sqrt{2})$$

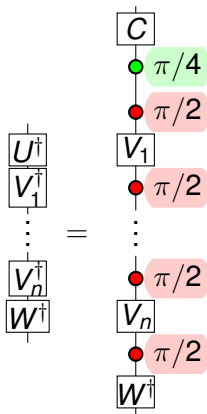
where $x_1, x_2, y_1, y_2, z_1, z_2 \in \mathbb{Z}$. By parity arguments, can show none of them represent $|0\rangle$, thus no normal form operator is equal to the identity.

The inverse of a normal form diagram has the same number of $\pi/4$ phase shifts as the original

We have $U^\dagger = \begin{array}{|c} \boxed{C} \\ \hline \bullet \pi/4 \end{array}$ and $V^\dagger = \begin{array}{|c} \bullet \pi/2 \\ \hline \boxed{V} \\ \hline \bullet \pi/2 \end{array}$, so

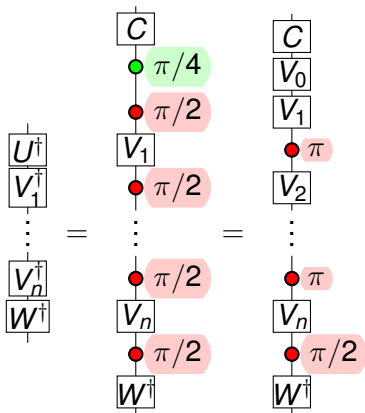
The inverse of a normal form diagram has the same number of $\pi/4$ phase shifts as the original

We have $U^\dagger = \begin{array}{|c} \boxed{C} \\ \bullet \pi/4 \end{array}$ and $V^\dagger = \begin{array}{|c} \bullet \pi/2 \\ \boxed{V} \\ \bullet \pi/2 \end{array}$, so



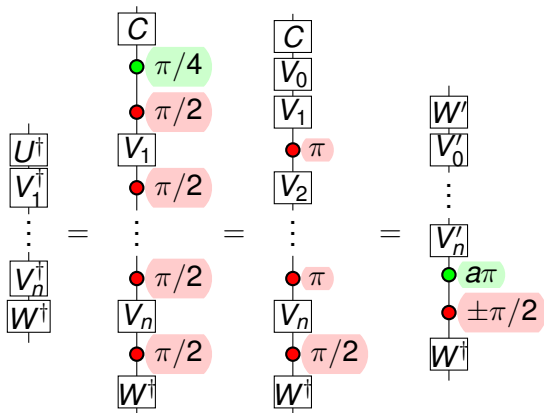
The inverse of a normal form diagram has the same number of $\pi/4$ phase shifts as the original

We have $U^\dagger = \begin{array}{c} \boxed{C} \\ \bullet \pi/4 \end{array}$ and $V^\dagger = \begin{array}{c} \bullet \pi/2 \\ \boxed{V} \\ \bullet \pi/2 \end{array}$, so



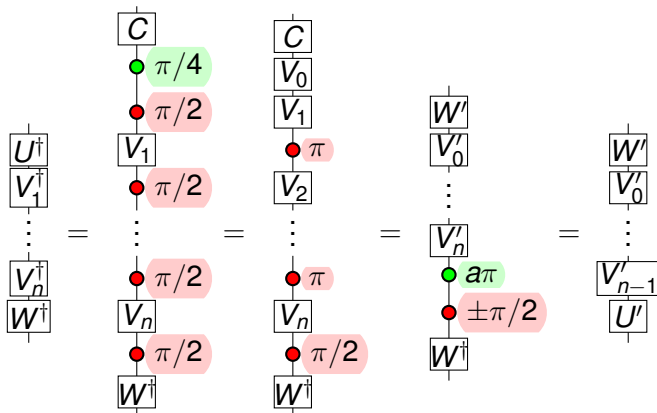
The inverse of a normal form diagram has the same number of $\pi/4$ phase shifts as the original

We have $U^\dagger = \begin{array}{c} \boxed{C} \\ \bullet \pi/4 \end{array}$ and $V^\dagger = \begin{array}{c} \bullet \pi/2 \\ \boxed{V} \\ \bullet \pi/2 \end{array}$, so



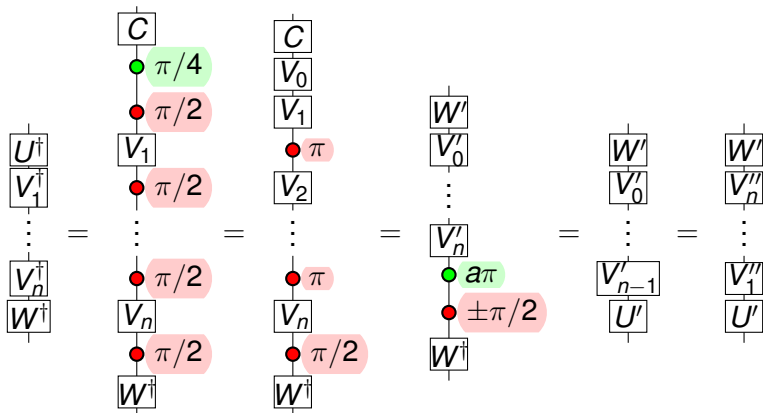
The inverse of a normal form diagram has the same number of $\pi/4$ phase shifts as the original

We have $U^\dagger = \begin{array}{|c|} \hline C \\ \hline \bullet \pi/4 \\ \hline \end{array}$ and $V^\dagger = \begin{array}{|c|} \hline \bullet \pi/2 \\ \hline V \\ \hline \bullet \pi/2 \\ \hline \end{array}$, so



The inverse of a normal form diagram has the same number of $\pi/4$ phase shifts as the original

We have $U^\dagger = \begin{array}{|c|} \hline C \\ \hline \bullet \pi/4 \\ \hline \end{array}$ and $V^\dagger = \begin{array}{|c|} \hline \bullet \pi/2 \\ \hline V \\ \hline \bullet \pi/2 \\ \hline \end{array}$, so



Normal forms are unique

Given two normal form diagrams that are not identical, composing one diagram with the inverse of the other will yield a non-trivial diagram.

- ▶ Assume diagrams are such that the topmost nodes differ (in colour or phase, or both).
- ▶ Inverting a diagram does not change the number of $\pi/4$ phases.
- ▶ Distinguish cases to see that the resulting diagram never collapses to the trivial one.

Outline

Motivation

Approximate completeness for single qubits

Completeness for stabilizer diagrams

Conclusions & Outlook

Stabilizer quantum mechanics

Stabilizer operations:

- ▶ preparation of qubits in state $|0\rangle$
- ▶ Clifford unitaries, generated by

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}, \quad H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad \Lambda X = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

- ▶ measurements in computational basis

Stabilizer quantum mechanics

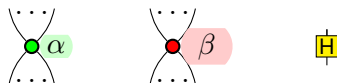
Stabilizer operations:

- ▶ preparation of qubits in state $|0\rangle$
- ▶ Clifford unitaries, generated by

$$S = \begin{pmatrix} 1 & 0 \\ 0 & i \end{pmatrix}, \quad H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}, \quad \Lambda X = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

- ▶ measurements in computational basis

ZX-calculus:



where α, β are multiples of $\pi/2$

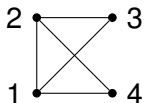
Graph states in the ZX-calculus

Definition

Let G be a finite simple undirected graph. The ZX-calculus diagram for the corresponding graph state consists of:

- ▶ for each node in G , a green node with one output, and
- ▶ for each edge in G , an edge with a Hadamard node on it.

E.g.



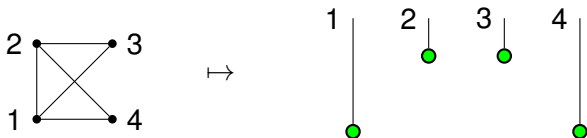
Graph states in the ZX-calculus

Definition

Let G be a finite simple undirected graph. The ZX-calculus diagram for the corresponding graph state consists of:

- ▶ for each node in G , a green node with one output, and
- ▶ for each edge in G , an edge with a Hadamard node on it.

E.g.



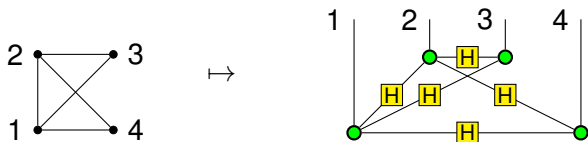
Graph states in the ZX-calculus

Definition

Let G be a finite simple undirected graph. The ZX-calculus diagram for the corresponding graph state consists of:

- ▶ for each node in G , a green node with one output, and
- ▶ for each edge in G , an edge with a Hadamard node on it.

E.g.



Local complementations

The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Local complementations

The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

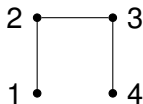
Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Theorem (Van den Nest et. al, 2004)

Two graph states are local Clifford-equivalent if and only if they are related by a sequence of local complementations.

A local complementation about a vertex v inverts the subgraph generated by the neighbourhood of v : e.g.



Local complementations

The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

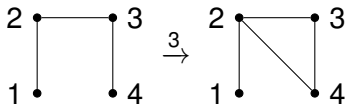
Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Theorem (Van den Nest et. al, 2004)

Two graph states are local Clifford-equivalent if and only if they are related by a sequence of local complementations.

A local complementation about a vertex v inverts the subgraph generated by the neighbourhood of v : e.g.



Local complementations

The local Clifford group consists of all tensor products of the single-qubit Clifford operators $\langle S, H \rangle$.

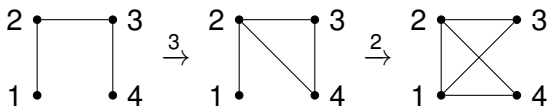
Theorem (Van den Nest et. al, 2004)

Any stabilizer state is local Clifford-equivalent to some graph state.

Theorem (Van den Nest et. al, 2004)

Two graph states are local Clifford-equivalent if and only if they are related by a sequence of local complementations.

A local complementation about a vertex v inverts the subgraph generated by the neighbourhood of v : e.g.

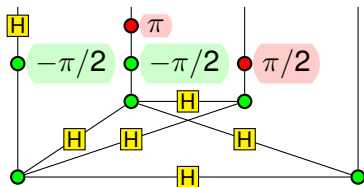


Stabilizer completeness proof (overview)

[Duncan & Perdrix, 2009] show that local complementations can be derived from the rules of the ZX-calculus.

Use this to show that the results from [Van den Nest et al., 2004] hold in the ZX-calculus:

- ▶ Any stabilizer ZX-calculus diagram can be rewritten into a graph state with local Clifford operators, e.g.



- ▶ There exists a terminating algorithm that, given two stabilizer diagrams, will rewrite them to be identical if possible.

See arXiv:1307.7025 for details.

Conclusions & Outlook

- ▶ ZX-calculus is not complete in general but fragments of it are complete, e.g.
 - ▶ line graphs where all phases are multiples of $\pi/4$ (single-qubit Clifford+T group)
 - ▶ diagrams where all phases are multiples of $\pi/2$ (stabilizer quantum mechanics)
- ▶ can these results be combined to multi-qubit Clifford+T operators?
- ▶ can we introduce some notion of approximate equality in the ZX-calculus?

Conclusions & Outlook

- ▶ ZX-calculus is not complete in general but fragments of it are complete, e.g.
 - ▶ line graphs where all phases are multiples of $\pi/4$ (single-qubit Clifford+T group)
 - ▶ diagrams where all phases are multiples of $\pi/2$ (stabilizer quantum mechanics)
- ▶ can these results be combined to multi-qubit Clifford+T operators?
- ▶ can we introduce some notion of approximate equality in the ZX-calculus?

Thank you!